

Nuevas Propiedades de la función Di-Bessel de Exton

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Resumen

En este trabajo se estudian nuevas propiedades de la función Di-Bessel de Exton A_μ . Esta función se define en forma similar a la función de Bessel de Primera clase y orden $\mu: J_\mu$. Como consecuencia, se demuestran propiedades del tipo de las conocidas para esta última función. Se obtienen entonces, reglas operacionales, ecuaciones diferenciales, relaciones de recurrencia, transformaciones integrales, representaciones integrales y desarrollos asintóticos.

New Properties of Di-Bessel function of Exton

In this paper we study certain new properties of the Exton's Di-Bessel function A_μ . The function A_μ is defined likeness to the Bessel function J_μ of the first kind and order μ . Then, properties like the ones which are verified by J_μ can be obtained for A_μ . Operational rules, differential equations, recurrence relations, integral transformations, integral representations and asymptotic expansions are thus obtained.

Key Words: Di-Bessel function, Recurrence formulae, Integral Representations, Asymptotic Behaviour.

1. Introduction

Exton [9] studied an extension of the Bessel function of the first kind an order μ by defining what he called a Di-Bessel function. In particular he defined the function $A_\mu(x)$:

$$A_\mu(x) = \sum_{m=0}^{\infty} \frac{(-1)^m (x/2)^{\mu+2m}}{(m!)^2 (\Gamma(\mu+1+m))^2} \quad (1.1)$$

This function $A_\mu(x)$, for x a complex variable and μ a fixed complex number is analytic on the x plane except possibly for branch points at $x = 0$ and $x = \infty$. When μ is an integer number we have the following generating function of the Di-Bessel function of integer order (Exton, 9)

$$I_0(2\sqrt{x}t) J_0(-2\sqrt{x}/t) = \sum_{n=-\infty}^{\infty} t^n A_n(x) \quad (1.2)$$

where J_0 is the Bessel function of the first kind and order 0, and I_0 denotes the modified Bessel function of the first kind and order 0.

Exton obtained a number of properties of this function including an orthogonal property. He showed that the differential equation satisfied by the Di-Bessel function is self-adjoint, being that some infinite integrals with the A_μ function were studied too.

The function $A_\mu(x)$ was generalized by A. Kumar Agarwal [3] by defining a n -ple Bessel function $A_{\mu,n}(x)$ which reduces for $n = 1$ and $n = 2$ to the Bessel and Di-Bessel functions, respectively. More recently, Sarabia [14] proves two recurrence relations for the function $A_\mu(x)$. He obtains an integral representation and also studies the oscillatory character of $A_\mu(x)$.

The object of this paper is to investigate $A_\mu(x)$ and study some of its new properties. Then,

as the Di-Bessel function appears as a particular case of the hyper-Bessel functions (Delerue, [4]; Dimovski-Kiryakova, [5]; Adamchik, [1] a good number of properties of the function A_μ can be proved. In particular, we obtain in this paper some differential equations which are verified for modified versions of the Di-Bessel function; we study two new recurrence relations for A_μ ; special definite integrals for A_μ are also obtained; an integral representation of the Di-Bessel function is likewise found. Lastly we present the asymptotic expansions of the A_μ function.

The present paper will be continued in a new work where we will define a new integral transformation of the Hankel type which contains the Di-Bessel function A_μ in its kernel. Then inversion formulae, convergence theorems, and spaces of test functions where this transformation will be a linear and continuous mapping will be studied.

2. Differential Equations

On investigating differential equations associated with the Di-Bessel function, we begin with the operational equation:

$$\delta^2 (\delta + \mu)^2 [E_\mu(\lambda x)] = \lambda x E_\mu(\lambda x) \quad (2.1)$$

where $\delta = x \frac{d}{dx}$ and $x > 0, \lambda > 0$ which is satisfied by the function:

$$E_\mu(\lambda x) = (-\lambda x)^{-\mu/2} A_\mu(2i\sqrt{\lambda x})$$

(A.K. Agarwal, [3], p. 142).

But, being that:

$\delta [x^\alpha f(x)] = x^\alpha (\delta + \alpha) f(x)$, for f a regular function and:

$$\delta [A_\mu(2i\sqrt{x})] = 2^{-1} [\delta (A_\mu)] (2i\sqrt{x})$$

we will find that:

$$\begin{aligned} \delta^2 (\delta + \mu)^2 (-x)^{-\mu/2} A_\mu(2i\sqrt{x}) &= \\ = (-1)^{-\mu/2} x^{-\mu/2} (\delta - \mu/2) (\delta - \mu/2) & \\ (\delta + \mu/2) (\delta + \mu/2) (A_\mu(2i\sqrt{x})) &= \end{aligned}$$

$$\begin{aligned} (-x)^{-\mu/2} \frac{1}{16} (\delta - \mu)^2 (\delta + \mu)^2 [A_\mu(u)] (2i\sqrt{x}) &= \\ (-x)^{-\mu/2} x A_\mu(2i\sqrt{x}) & \end{aligned}$$

In this way we have proved that the function $A_\mu(x)$ satisfies the operational equation:

$$(\delta - \mu)^2 (\delta + \mu)^2 [A_\mu(x)] = -4 x^2 A_\mu(x) \quad (2.2)$$

When we choose the function $B_{\alpha,\beta,\mu}(x) = x^\alpha A_\mu(Kx^\beta)$ for a constant $K, K > 0$, we will find:

$$A_\mu(y) = (y/K)^{-\alpha/\beta} B_{\alpha,\beta,\mu}(y/K)^{1/\beta} ; y > 0$$

and being that:

$$\delta [f((y/K)^{1/\beta})] = (\beta^{-1} \delta (f)) ((y/K)^{1/\beta})$$

for f a regular function; we will obtain:

$$\begin{aligned} (\delta - \mu)^2 (\delta + \mu)^2 [(y/K)^{-\alpha/\beta} B_{\alpha,\beta,\mu}((y/K)^{1/\beta})] &= \\ -4y^2 [(y/K)^{-\alpha/\beta} B_{\alpha,\beta,\mu}((y/K)^{1/\beta})] & \end{aligned}$$

and so we can find the operational relation:

$$\begin{aligned} (\delta - \mu\beta - \alpha)^2 (\delta + \mu\beta - \alpha)^2 [B_{\alpha,\beta,\mu}(y)] &= \\ -y^{2\beta} 4 K^2 B_{\alpha,\beta,\mu}(y) \beta^4 & \end{aligned} \quad (2.3)$$

In particular, when we choose $\alpha = 1/2$ and $\beta = 1$, we obtain:

$$\begin{aligned} (\delta - \mu - 1/2)^2 (\delta + \mu - 1/2)^2 B_{1/2,1,\mu}(y) &= \\ -y^2 4 K^2 B_{1/2,1,\mu}(y) & \end{aligned} \quad (2.3)$$

and as:

$$(\delta + \gamma) (\delta + \rho) = \delta^2 + (\gamma + \rho) \delta + \gamma \rho, \text{ for } \gamma, \rho \in \mathbb{R}$$

thus:

$$\begin{aligned} [\delta^4 - 2\delta^3 - (2\mu^2 - 3/2)\delta^2 + 2(\mu - 1/2)(\mu + 1/2)\delta + \\ + (\mu + 1/2)^2 (\mu - 1/2)^2] [\sqrt{x} A_\mu(Kx)] &= \\ -x^2 [4 K^2 A_\mu(Kx)] \sqrt{x} & \end{aligned}$$

i.e.:

$$\left[x^2 \frac{d^2}{dx^2} - (\mu^2 - 1/4) \right] \left[x^2 \frac{d^2}{dx^2} - (\mu^2 - 1/4) \right] [\sqrt{x} A_\mu(Kx)] = -x^2 4 K \sqrt{x} A_\mu(x)$$

Furthermore, if $S_\mu = \frac{d^2}{dx^2} \cdot \frac{4\mu^2 - 1}{4x^2}$ it denotes the Bessel's classical differential operator, we have found the following formula (2.4):

$$[S_\mu X^2 S_\mu \sqrt{x}] [A_\mu(Kx)] = -4 K^2 (\sqrt{x} A_\mu(Kx))$$

3. Recurrence Formulae For $A_\mu(x)$

In this section we obtain certain new recurrence relations for the Di-Bessel function $A_\mu(x)$. In Sarabia 14 we can find the following recurrence formulae for this function:

1. $\frac{d}{dx} x \frac{d}{dx} (x^\mu A_\mu(x)) = 2 x^\mu A_{\mu-1}(x)$
2. $\frac{d}{dx} x \frac{d}{dx} (x^{-\mu} A_\mu(x)) = -2 x^{-\mu} A_{\mu+1}(x)$
3. $2_\mu A'_\mu(x) = A_{\mu-1}(x) + A_{\mu+1}(x)$
4. $\mu^2 A_\mu(x) + (1 - 2\mu) x A'_\mu(x) + x^2 A''_\mu(x) = -2x A_{\mu+1}(x)$
5. $\mu^2 A_\mu(x) + (1 + 2\mu) x A'_\mu(x) + x^2 A''_\mu(x) = 2x A_{\mu-1}(x)$
6. $(\mu-1)x^2 A_{\mu+2}(x) + 2(\mu-1)(\mu+1)(2\mu+1)x A_{\mu+1}(x) + 2_\mu(2\mu^2(\mu-1)(\mu+1) + x^2) A_\mu(x) + 2(\mu-1)(\mu+1)(1-2\mu)x A_{\mu-1}(x) + (\mu+1)x A_{\mu-2}(x) = 0$

Additionally, if we define the differential operators N_μ and M_μ by:

$$N_\mu = x^{\mu+1/2} \frac{d}{dx} x \frac{d}{dx} x^{-\mu-1/2}$$

$$M_\mu = x^{-\mu-1/2} \frac{d}{dx} x \frac{d}{dx} x^{\mu+1/2}$$

we find that:

$$x^2 M_\mu N_\mu = x^2 (x^{\mu-1/2} \frac{d}{dx} x \frac{d}{dx} x^{\mu+1/2})$$

$$(x^{\mu+1/2} \frac{d}{dx} x \frac{d}{dx} x^{-\mu-1/2}) = S_\mu X^2 S_\mu$$

and then:

$$N_\mu [\sqrt{x} A_\mu(x)] = -2 \sqrt{x} A_{\mu+1}(x)$$

$$M_\mu [\sqrt{x} A_{\mu+1}(x)] = 2 \sqrt{x} A_\mu(x)$$

Also, if R_μ denotes the Riemann-Liouville operator (Ederlyi, [8]) given in (3.1):

$$R_\mu f(x) = \Gamma(\mu)^{-1} \int_0^x f(t) (x-t)^{\mu-1} dt, \tag{3.1}$$

for a regular function $f, \mu > 0, x > 0,$

we know that:

$$\frac{x^{\mu+n}}{\Gamma(\mu+n)} = \frac{1}{\Gamma(n+1)} R_\mu [x^n] \quad n \in \mathbb{N}, \text{ and } \mu > 0$$

Then, being as:

$$[(x/2)^\mu A_\mu(x)] (\sqrt{x}) = \sum_{n=0}^{\infty} \frac{(-1)^n (x/2)^{\mu+n}}{(n!)^2 [\Gamma(\mu+n+1)]^2} (1/2)^{\mu+1}$$

we have:

$$[(x/2)^\mu A_\mu(x)] (\sqrt{x}) = \sum_{n=0}^{\infty} \frac{R_\mu [(-1)^n (x/2)^n]}{(n!)^3 \Gamma(\mu+n+1)} (1/2)^{2\mu+n}$$

$$= R_\mu \left[(x/2)^{-\mu} \sum_{n=0}^{\infty} \frac{(-1)^n (x/2)^{n+\mu}}{(n!)^3 \Gamma(\mu+n+1)} (1/2)^{2\mu+n}(x) \right] =$$

$$= R_\mu \left[x^{-\mu} \sum_{n=0}^{\infty} \frac{R_\mu [(-1)^n (x/2)^n]}{(n!)^2} (1/2)^{2\mu+n} \right] (x) =$$

$$R_\mu \left[x^{-\mu} R_\mu \sum_{n=0}^{\infty} \frac{(-1)^n (x/2)^n}{(n!)^4} (1/2)^n \frac{1}{2^{2\mu}} \right] (x) =$$

$$= [R_\mu x^{-\mu} R_\mu] [A_0(\sqrt{x})] (x) 2^{-2\mu}$$

Then we will have obtained the general formulae that follows:

$$A_\mu(x) = (2x)^\mu R_\mu x^{-\mu} R_\mu [A_0(\sqrt{x})] (x^2) \tag{3.2}$$

4. Infinite Integrals

Consider, first of all, the integral

$$I = \int_0^{\infty} e^{-ax} x^{v-1} A_\mu(bx) dx \tag{4.1}$$

$Re a > 0, Re(v+\mu) > 0$

Re $a > 0, Re(v+\mu) > 0$

In Exton [9], p. 860 it has been found that "I" can be expressed in the form:

$$a^{-v} (b/2a)^\mu \frac{\Gamma(\mu+v)}{[\Gamma(\mu+1)]^2} {}_2F_3((v+\mu)/2, (v+\mu+1)/2; \mu+1, \mu+1, 1; -b^2/a^2) \tag{4.2}$$

Sarabia, [14], p. 72 obtained an other defined integral for $A_\mu(x)$; but it is not of very great interest for usm. We are interested in the application of some integral transformations to $A_\mu(x)$. So, in (4.1)-(4.2) we can find the Laplace transformation of A_μ . Then this transformation $L(A_\mu(x))(a)$ is given in (4.3):

$$L(A_\mu(x))(a) = 2^{-\mu} a^{-\mu-1} \sum_{n=0}^{\infty} \frac{\Gamma(\mu+1+2n)}{(n!)^2 [\Gamma(\mu+1+n)]^2} (-1/a^2)^n \tag{4.3}$$

Re $a > 0$

Moreover, we know the hyper-Bessel functions defined by Delerue, [4] in 1953 in the form:

$$\frac{x^{\lambda_1+\lambda_2+\dots+\lambda_n}}{n+1} \sum_{n=0}^{\infty} \frac{J_{\lambda_1, \lambda_2, \dots, \lambda_n}^{(n)}(x)}{\Gamma(\lambda_1+1+n) \dots \Gamma(\lambda_n+1+n)} \tag{4.4}$$

$\lambda_1, \lambda_2, \dots, \lambda_n \in \mathbb{R}, n \in \mathbb{N}, x > 0$

and being that $A_\mu(x)$ coincides with the function:

$$A_\mu(x) = J_{0, \mu, \mu}^{(3)}(2\sqrt{2x})$$

we find (Delerue, [4], p. 233):

$$L \left[x^{(n\lambda_1-\lambda_2-\dots-\lambda_n)/(n+1)} J_{\lambda_1, \lambda_2, \dots, \lambda_n}^{(n)}(n+1)^{n+1} \sqrt{x} \right] (a) = a^{(n\lambda_1-\lambda_2-\dots-\lambda_n)/(n+1)} J_{\lambda_1, \lambda_2, \dots, \lambda_n}^{(n)}(n/\sqrt{a}) \tag{4.6}$$

and so:

$$L \left[x^{-\mu/2} A_\mu(\sqrt{x}/4) \right] (a) = a^{\mu/3} J_{\mu, \mu}^{(2)}(3a^{-1/3}) \tag{4.7}$$

The Mellin transformation $M(A_\mu)(s)$ of A_μ will be given in the form:

$$M(A_\mu(x))(s) = \int_0^\infty x^{s-1} A_\mu(x) dx \tag{4.8}$$

and we can obtain an explicit expression of (4.8) when the integral representation of A_μ is known which will be studied in the paragraph 5...

5. Integral Representations of the Di-Bessel Function

We shall next examine various representations of the Di-Bessel function by a system of definite integrals and contour type integrals. Firstly we find the following integral representation obtained by Sarabia, [14], p. 67:

$$A_\mu(x) = \frac{x^\mu}{2\pi i \Gamma(\mu+1) 2^\mu} \int_C e^s s^{-(\mu+1)} {}_0F_2(-; \mu+1; -x^2/4s) ds \tag{5.1}$$

where C denotes Hankel's contour.

When μ is a nonnegative integer $\mu = n \in \mathbb{N}$, the expression of (5.1) will stay in the form (Exton, [9], p. 857):

$$A_n(x) = \frac{1}{2\pi i} \int_{-\infty}^{(0)^+} S^{-n-1} I_0(2\sqrt{XS}) J_0(-2\sqrt{XS}) ds \tag{5.2}$$

where I_0 and J_0 denote the classical Bessel functions of order 0.

We now take into account Hankel's well-known integral representation of Bessel functions, and then:

$$\begin{aligned} \pi [\Gamma(\mu+1/2)]^2 2^\mu \sum_{n=0}^{\infty} \frac{(-1)^n (x/2)^{\mu+2n}}{n!^2 [\Gamma(\mu+n+1)]^2} &= \\ \sum_{n=0}^{\infty} \frac{(-1)^n 2^{2n} x^{\mu+2n}}{[(2n)!]^2} \left[\frac{\Gamma(n+1/2) \Gamma(\mu+1/2)}{\Gamma(n+\mu+1)} \right]^2 &= \\ = \sum_{n=0}^{\infty} \frac{2^{2n} x^{\mu+2n} (-1)^n}{(2n!)^2} \int_0^1 t^{2n} &= \\ (t^2-1)^{\mu-1/2} dt \int_0^1 \tau^{2n} (\tau^2-1)^{\mu-1/2} d\tau &= \\ = \sum_{n=0}^{\infty} \frac{(2n!)^2}{(n!)^2} \frac{x^{\mu+n} i^n}{(n!)^2} \int_{-1}^1 t^{2n} &= \\ (t^2-1)^{\mu-1/2} dt \int_{-1}^1 \tau^n (\tau^2-1)^{\mu-1/2} d\tau & \end{aligned}$$

Consider the function obtained by interchanging the signs of summation and integration on the right; it is:

$$C_0(-2ixt\tau) (t^2-1)^{\mu-1/2} (\tau^2-1)^{\mu-1/2}$$

where:

$$C_0(z) = \sum_{n=0}^{\infty} \frac{(-z)^n}{[n!]^2}$$

denotes the Bessel-Clifford function of the first kind and order 0 (Hayek, [11])

This is an analytic of x for all values of x , and when expanded in ascending powers of x by MacLaurin's theorem, the coefficients may be obtained by differentiating with regard to x under the double integral sing and making x zero after differentiation. Hence:

$$x^\mu \int_{-1}^1 \int_{-1}^1 C_0(-2ixt\tau) (t^2-1)^{\mu-1/2} (\tau^2-1)^{\mu-1/2} dt d\tau = \pi [\Gamma(\mu+1/2)]^2 2^\mu \sum_{n=0}^\infty \frac{(-1)^n (x/2)^{\mu+2n}}{(n!)^2 (\Gamma(\mu+1+n))^2};$$

i.e., the function $A_\mu(x)$ admits the following integral representation (5.3):

$$A_\mu(x) = \frac{(x/2)^\mu}{\pi [\Gamma(\mu+1/2)]^2} \int_{-1}^1 \int_{-1}^1 C_0(-2ixt\tau) (t^2-1)^{\mu-1/2} (\tau^2-1)^{\mu-1/2} dt d\tau \tag{5.3}$$

Furthermore, by the formula (4.6) we obtain:

$$L[x^{\mu/2} J_{\mu,\mu,0}^{(3)}(4\sqrt{x})] (a) = a^{-2\mu/3} J_{\mu,0}^{(2)}(3/\sqrt[3]{a})$$

and then, L^{-1} being the inverse Laplace transformation; we find:

$$L^{-1} [a^{-2\mu/3} J_{\mu,0}^{(2)}(3/\sqrt[3]{a})] (a+x) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} e^{ax} a^{-2\mu/3} \sum_{n=0}^\infty \frac{(-1)^n 3^n a^{-n/3}}{(n!)^2 \Gamma(\mu+n+1)} da$$

for a real number $c \in \mathbb{R}$.

So, if we take Hankel's well-known generalisation of the second Eulerian integral:

$$\frac{1}{\Gamma(\mu+m+1)} = \frac{1}{2\pi i} \int_{-\infty}^{(0^+)} t^{\mu-m-1} e^t dt$$

in which the phase of t increases from $-\pi$ to π as t describes the contour, we could carry out calculations analogous to the previous ones and in that way obtain:

$$x^{\mu/2} J_{\mu,\mu,0}^{(3)}(4\sqrt{x}) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} e^{ax} a^{-2\mu/3}$$

$$\frac{1}{2\pi i} \int_{-\infty}^{(0^+)} C_0(3/\sqrt[3]{a}) e^{t} t^{-\mu-1} dt da$$

$$Re a > 0$$

i.e.; A_μ will allow the new integral representation:

$$A_\mu(x) = x^{-\mu/2} \frac{1}{2\pi i} \int_{-\infty}^{(0^+)} e^t t^{-\mu-1} L^{-1} [a^{-2\mu/3} C_0(3a^{-1/3}t)] (x^2/4) dt \tag{5.4}$$

Lastly, we can use the formula (5.3) to obtain the Mellin transformation of $A_\mu(x)$

$$M(A_\mu(x))(s) = \int_0^\infty x^{s-1} \left[\int_{-1}^1 (t^2-1)^{\mu-1/2} \left[\int_{-1}^1 (\tau^2-1)^{\mu-1/2} C_0(-2ixt\tau) dt \right] dt \right] dx \cdot (x/2)^\mu \pi^{-1} \Gamma(\mu+1/2)^{-2},$$

and interchanging the integration symbols, we shall obtain:

$$M(A_\mu(x))(s) = \int_{-1}^1 (t^2-1)^{\mu-1/2} \left[\int_{-1}^1 (\tau^2-1)^{\mu-1/2} \frac{2^{-\mu}}{\pi [\Gamma(\mu+1/2)]^2} \cdot \left[\int_0^\infty x^{\mu+s-1} J_0(-2\sqrt{2ixt\tau}) dx \right] dt \right] dt = \frac{2^{-\mu+1} 2^{2\mu+2s-1} \Gamma(\mu+s)}{\pi \Gamma(\mu+1/2)^2 \Gamma(1-\mu-s)}$$

$$\left[\int_{-1}^1 (t^2-1)^{\mu-1/2} t^{-\mu-s} dt \right]^2 2^{-3(\mu+s)} 1^{-(\mu+s)} =$$

for $3/2 > Re s > 0$, according to Ditkin-Proudnikov, [7], p. 359

So:

$$M(A_\mu(x))(s) = \frac{2^{-2\mu-s} i^{-\mu-s} \Gamma(\mu+s)}{\pi [\Gamma(\mu+1/2)]^2 \Gamma(1-\mu-s)} \beta(\mu+1/2, 1/2-\mu/2-s/2) (1+(1)^{\mu+s})/2 \tag{5.5}$$

Then, we shall obtain:

$$M(A_\mu(x))(s) = \pi^{-1} 2^{-2-2\mu-s} (-1)^{-(\mu+s)/2} (1+(-1)^{\mu+s}) \frac{\Gamma(\mu+s) \Gamma((1-\mu-s)/2)}{\Gamma(1-\mu-s) \Gamma(1+\mu/2-s/2)} \cdot \frac{\Gamma(1/2-\mu/2-s/2)}{\Gamma(1+\mu/2-s/2)}$$

for $3/2 > Re s > 0$ and $\mu > -1/2$.

Moreover, the following relations are known:

$$[1 + (-1)^{\mu-s}]^2 = 2(1 + (-1)^{\mu+s})$$

and

$$(-1)^{-(\mu+s)/2} [1 + (-1)^{\mu+s}]^2 = 4 \cos(\mu+s)\pi/2$$

then:

$$M(A_\mu(x))(s) = \pi^{-1} 2^{-s-2\mu} \cos(\mu+s)\pi/2 \frac{\Gamma(\mu+s)}{\Gamma(1-\mu-s)} \left[\frac{\Gamma(1/2 - \mu/2 - s/2)}{\Gamma(1 + \mu/2 - s/2)} \right]^2$$

but, by using the well known formulae:

$$\sin(\pi z) \Gamma(z) \Gamma(1-z) = \pi$$

and

$$\sqrt{\pi} \Gamma(2z) = 2^{2z-1} \Gamma(z) \Gamma(z+1/2)$$

we can operate on (5.5) until we obtain:

$$M(A_\mu(x))(s) = \pi^{-1} 2^{\mu+s-1} \sin \pi(\mu+s)/2 \left[\frac{\Gamma(\mu/2 + s/2)}{\Gamma(1 + \mu/2 - s/2)} \right]^2 \quad (5.6)$$

6. The Asymptotic Behaviour of A_μ

In Mathai - Saxena [12], p. 43 we find that:

$G_{p,q}^{m,n} \left[z / \begin{matrix} a_1, a_2, \dots, a_p \\ b_1, b_2, \dots, b_q \end{matrix} \right]$ being Meijer's G-function which was defined by Meijer in 1936 by means of a finite series of the generalized hypergeometric function; then the generalized hypergeometric function ${}_pF_q(a_1, a_2, \dots, a_p; b_1, b_2, \dots, b_q; z)$ admits the following representation through the function G; i. e.:

$${}_pF_q(a_1, a_2, \dots, a_p; b_1, b_2, \dots, b_q; z) = \frac{\prod_{j=1}^q \Gamma(b_j)}{\prod_{j=1}^p \Gamma(a_j)} G_{p,q+1}^{1,p} \left[z / \begin{matrix} 1-a_1, \dots, 1-a_p \\ 0, 1-b_1, \dots, 1-b_q \end{matrix} \right]$$

when certain conditions are satisfied. Then, our function $A_\mu(x)$ can be expressed in the form:

$$[z^{-\mu} A_\mu(z)](2\sqrt{z}) = 2^{-\mu} F_3(-; 1, \mu+1, \mu+1; -z) = 2^{-\mu} \Gamma(\mu+1)^2 G_{0,4}^{1,0} \left[z / \begin{matrix} - \\ 0, 0, -\mu, -\mu \end{matrix} \right] \quad (6.1)$$

(6.1)

Then, it is possible to use the theorems about the asymptotic expansions of Meijer's function in Mathai-Saxena [12]; § 1.8. to obtain that:

" there exist constants: $M_1, M_2, \dots, M_n, \dots$ and:

$$G_{0,4}^{1,0} \left[z / \begin{matrix} - \\ 0, 0, -\mu, -\mu \end{matrix} \right] = \sum_{i=1}^{\infty} (-2\pi i)^{-3} e^{2\mu\pi i} \exp -4 (ze^{i\pi 3})^{1/4}$$

$$\cdot (ze)^{-(4\mu+3)/8} \left[\frac{(2\pi)^{3/2}}{2} + \frac{M_1}{(ze^{i\pi 3})^{1/4}} + \frac{M_2}{(ze^{i\pi 3})^{1/2}} + \dots \right]$$

In this way, we can find the following asymptotic expansion of the Di-Bessel function A_μ :

$$A_\mu(z) = \Gamma(\mu+1)^2 (-1)^{(4\mu+3)/8} (2\pi)^{-3} 2^{7/4} \underset{arg z = 0}{z \rightarrow \infty} \sin 2\mu\pi e^{-2\sqrt{2z}i} z^{-3/4}$$

$$\left[\frac{(2\pi)^{3/2}}{2} + \frac{M_1}{z^{1/2}} + \frac{M_2}{z} + \dots \right] \quad (6.2)$$

where $M_1, M_2, \dots, M_n, \dots$ are appropriate constants

The asymptotic behaviour of A_μ can be obtained from (6.2) or seeing that, according to Mathai-Saxena, [12], p. 307, the growth of $G_{p,q}^{m,0}$ in $z \rightarrow \infty$ is given in:

$$G_{p,q}^{m,0} \left[z / \begin{matrix} a_1, a_2, \dots, a_p \\ b_1, b_2, \dots, b_q \end{matrix} \right] = O \left[\exp \left[(p-q) z^{1/(q-p)} \right] z^\beta \right] \quad (6.3)$$

where:

$$\beta = \frac{1}{q-p} \left[\frac{p-q+1}{2} + \sum_{i=1}^q b_i - \sum_{i=1}^p a_i \right]$$

$$|arg z| < (q-p+1)\pi$$

Thus, we shall find that:

$$A_\mu(x) = O \left[e^{-2\sqrt{2x}} x^{-3/4} \right] \underset{x \rightarrow \infty}{} \quad (6.4)$$

In the same form, it would easily be seen that:

$$A_\mu(x) = O[x^\mu] \underset{x \rightarrow 0}{} \quad (6.5)$$

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Recibido: 20 de Junio de 1991

En forma revisada: 11 de Febrero de 1992